

$$\begin{aligned}
\alpha & \frac{3\delta v_u}{B\delta 1} \frac{v\phi}{2v\phi\delta 1} \rho v_u \phi & ; & \quad \beta \frac{1}{4} \frac{\delta 1}{\delta 1} \frac{2v_u\phi\delta 1}{2v\phi\delta 1} \frac{v\phi}{v_u\phi} & ; & \quad \gamma \frac{1}{4} \frac{2GB\delta 1}{3\delta 1} \frac{v\phi\delta 1}{2v\phi\delta 1} \rho v_u \phi \\
\text{c} & \frac{1}{4} \frac{2GB^2\delta 1}{9\delta v_u} \frac{v\phi\delta 1}{v\phi\delta 1} \rho v_u \phi^2 k & ; & \quad \eta \frac{1}{4} \frac{\delta 1}{\delta 1} \frac{v\phi}{2v\phi} & ; & \quad \Theta \frac{1}{4} \frac{u_r}{r} \rho \frac{u_r}{r} \rho \frac{u_z}{z}
\end{aligned}
\tag{4}$$

In Eqn (1)–(4), G is the linear shear modulus and v is Poisson's ratio of the material (i.e. the undrained elastic compressibility); v_u is the undrained Poisson's ratio of the fluid-saturated medium; k is the hydraulic conductivity; B is Skempton's pore pressure coefficient [32]; and ρ is the mixture density of the lattice, $\rho = \rho_s \phi + \rho_f (1 - \phi)$.

$$\rho = \rho_s \phi + \rho_f (1 - \phi) \tag{5}$$

Certain thermodynamic constraints need to be satisfied, namely the density of the fluid ρ_f is constant [33]; it can be shown that the equilibrium condition can be expressed in the form: $G > 0$; $0 < B < 1$; $1 < v < v_u < 0.5$. Also note that $\rho = \rho_s \phi + \rho_f (1 - \phi)$ (see also [32] for details).

The accuracy of the eigenvalues in terms of $S(r, z, t)$ and $E(r, z, t)$ can be determined

$$\sigma_{zz} \partial r, 0, t \neq 0 \quad a < r < \infty \quad (16)$$

$$\sigma_{rz} \partial r, 0, t \neq 0 \quad 0 < r < \infty \quad (17)$$

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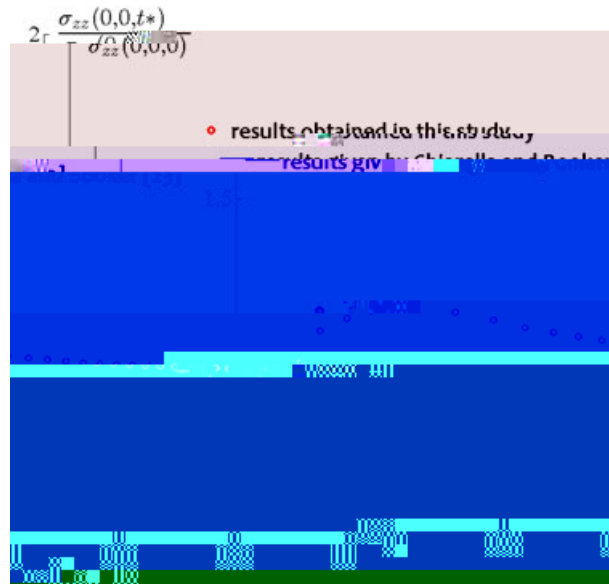
Because the concentration is related to the chemical face displacement, the efficiency
mainly, the filling rate is a linear function of the global equilibrium:

$$\sum_{i=1}^n \tilde{\sigma}_i$$

the maximum decrease of 7.6% when $\nu=0$, $\nu_u=0.5$ but decrease of 0.4% when $\nu=0.2$, $\nu_u=0.5$ and $\nu=0.4$, $\nu_u=0.5$.

Figure 3 compares the normalized vertical stress, $\sigma_{zz}(0,0,t^*)/\sigma_{zz}(0,0,0)$, in Chiaella and Bekke [23]. The trend of the normalized vertical stress in Chiaella and Bekke [23]. The decrease in the maximum normalized vertical stress is 3.5%. The decrease at $t^*=1$ is 1.6%.

Figure 4 compares the results obtained for Case I (completely pervious surface) and Case II (completely impervious surface). The condition of the increase in the vertical stress is shown in Figure 4.



been established in a similar manner. The corresponding boundary conditions are given by Yeh and Selvadurai [25] as follows: when $v_u = 0$, the boundary conditions are $\sigma_{zz} = \tau_{zx} = \tau_{zy} = 0$; when $v_u = 0.5$, the boundary conditions are $\sigma_{zz} = \tau_{zx} = \tau_{zy} = 0$, $\sigma_{zz} = 0$. The corresponding boundary conditions are given by Yeh and Selvadurai [25].

Figure 5 shows the effect of ν_u on the consolidation of Case I. When $\nu_u = 0$, the consolidation decreases as ν_u increases. On the other hand, when ν_u is fixed at 0.5, the consolidation increases as ν increases (see also Figure 4). The corresponding boundary conditions are given by Yeh and Selvadurai [25] as denoted by circles in Figure 5. The agreement is good (less than 4% when $t^* > 10^{-2}$) except for the case, $\nu = 0, \nu_u = 0.5$. The discrepancy becomes larger near the initial and final stages (see Table II).

The effect of ν_u on the consolidation of Case II is shown in Figure 6. As shown in Figure 5, the consolidation decreases as ν_u decreases when ν is fixed at 0, while the consolidation increases as ν increases when $\nu_u = 0.5$ (see also Figure 4). The corresponding boundary conditions are given by Yeh and Selvadurai [25] as denoted by circles in Figure 6. The agreement is

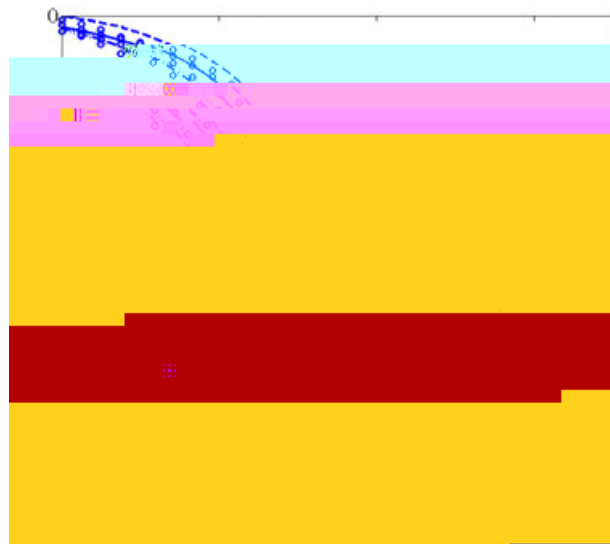


Figure 5. Consolidation of different ν_u when the surface is completely pervious.

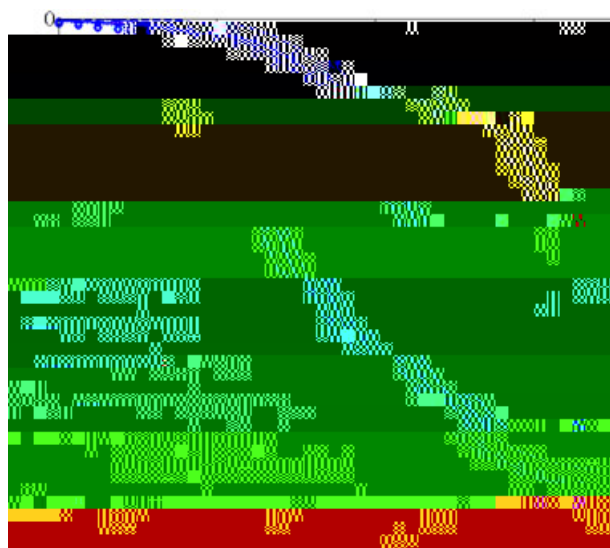


Figure 6. Consolidation of different ν_u when the surface is completely impervious.

The eigenvalues and the eigenfunctions given by Y. e and Sel ad ai [25] i again g d i h ma im me
of 6% e ce f the ca e $v=0$, $v_u=0.5$.

7. CONCLUDING REMARKS

Conclusion of the elastic half space in the semi ed b nda c ndi n a he lane face
of the half space in e m f the di lacemen , e e and e id e e. The fin ced e
of the semi ed b nda c ndi n in le a e f d al in eg al e a i n in he La lace an f m
d main ha cann be led ing a c n eni nal in eg al an f m (Hankel and La lace
an f m) a ach. An ale na i e a ach, he e he c n ac e i di c e i ed in

APPENDIX B

The natural eigenfunctions in the displacement in the z-direction are given by

Case I: The end face is completely permeable

$$u_z(r, z, t) = \frac{p^* a^*}{2G} \int_0^\infty \int_0^\infty \left\{ \frac{[\eta(\phi^2 - \xi^2) \delta \Gamma + \xi z \delta \Gamma - 1 + \delta \Gamma + 1 + \phi \delta 2 \Gamma \eta \xi - \xi^2] e^{-\xi z}}{\delta \phi - \xi^2 \eta \delta \Gamma - 1 + \delta \phi + \xi^2 \delta 2 \Gamma \eta \xi - \xi^2} \right.$$

$$\left. - \frac{\phi^2}{\phi^2} \right\}$$

17. Selvadurai APS (Ed).